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## A note on the product

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# A NOTE ON THE PRODUCT* $r^{-k} \cdot \nabla\left(\Delta r^{2-m}\right)$ 

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The difficulties inherent in defining products, powers or nonlinear operations of generalized functions have not prevented their appearance in the literature (see e.g. [1, 2, 3]). In [4], a definition for a product of distributions is given using a $\delta$-sequence. However, they show that $\delta^{2}$ does not exist. Extending definitions of products from one-dimensional space to $m$-dimensional by using appropriate delta-sequences has recently been an interesting topic in distribution theory. From a single variable function $\rho(s)$ (see [5]) defined on $R^{+}, \mathrm{Li}$ and Fisher introduce a 'workable' delta sequence in $R^{m}$ to deduce a non-commutative neutrix product of $r^{-k}$ and $\Delta \delta$ ( $\Delta$ denotes the Laplacian) for any positive integer $k$ between 1 and $m-1$ inclusive. In [6], Li provides a modified $\delta$-sequence and defines a 'bridge' distribution $\frac{d^{k}}{d r^{\delta}} \delta(x)$, which can be used to obtain the more general product of $r^{-k}$ and $\Delta^{\prime} \delta$. The object of this paper is to utilize the fact that $E=\frac{\Gamma(m / 2)}{(2-m) 2 \pi^{m / 2}} r^{2-m}$ is an elementary solution of Laplacian equation $\Delta u(x)=\delta(x)$ and to apply a much simpler version of the main result in [7] to derive the product $r^{-k} \cdot \nabla\left(\Delta r^{2-m}\right)$ where $m$ is any dimension greater than 2. Finally, we also compute the products $r^{-k} \cdot \Delta \ln r$ and $r^{-k} \cdot \nabla(\Delta \ln r)$ based on the equation $\frac{1}{2 \pi} \Delta \ln r=\delta$ for $m=2$.

KEY WORDS: Pizetti's formula, delta sequence, neutrix limit, neutrix product and distribution

MSC (2000): 46F10

## 1. THE DISTRIBUTION $r^{\boldsymbol{\lambda}}$

Let $r=\left(x_{1}^{2}+\ldots+x_{m}^{2}\right)^{1 / 2}$ and consider the functional $r^{\lambda}$ (see [8]) defined by

$$
\begin{equation*}
\left(r^{\lambda}, \phi\right)=\int_{R^{m}} r^{\lambda} \phi(x) d x \tag{1}
\end{equation*}
$$

[^0]where $\operatorname{Re} \lambda>-m$ and $\phi(x) \in \mathcal{D}_{m}$, the space of infinitely differentiable functions of the variable $x=\left(x_{1}, x_{2}, \ldots, x_{m}\right)$ with compact support. Because the derivative
$$
\frac{\partial}{\partial \lambda}\left(r^{\lambda}, \phi\right)=\int r^{\lambda} \ln r \phi(x) d x
$$
exists, the functional $r^{\lambda}$ is an analytic function of $\lambda$ for $\operatorname{Re} \lambda>-m$.
For $\operatorname{Re} \lambda \leq-m$, we should use the following identity (2) to define its analytic continuation. For $\operatorname{Re} \lambda>0$, we could deduce
$$
\Delta\left(r^{\lambda+2}\right)=(\lambda+2)(\lambda+m) r^{\lambda}
$$
simply by calculating the left-hand side. By iteration we find for any integer $k$ that
\[

$$
\begin{equation*}
r^{\lambda}=\frac{\Delta^{k} r^{\lambda+2 k}}{(\lambda+2) \cdots(\lambda+2 k)(\lambda+m) \cdots(\lambda+m+2 k-2)} \tag{2}
\end{equation*}
$$

\]

On making substitution of spherical coordinates in (1), we come to

$$
\begin{equation*}
\left(r^{\lambda}, \phi\right)=\int_{0}^{\infty} r^{\lambda}\left\{\int_{r=1} \phi(r \omega) d \omega\right\} r^{m-1} d r \tag{3}
\end{equation*}
$$

where $d \omega$ is the hypersurface element on the unit sphere. The integral appearing in the above integrand can be written in the form

$$
\begin{equation*}
\int_{r=1} \phi(r \omega) d \omega=\Omega_{m} S_{\phi}(r) \tag{4}
\end{equation*}
$$

where $\Omega_{m}=\frac{2 \pi^{m / 2}}{\Gamma(m / 2)}$ is the hypersurface area of the unit sphere imbedded in Euclidean space of $m$ dimensions, and $S_{\phi}$ is the mean value of $\phi$ on the sphere of radius $r$.

It was proved in [8] that $S_{\phi}(r)$ is infinitely differentiable for $r \geq 0$ and has bounded support, and that

$$
S_{\phi}(r)=\phi(0)+a_{1} r^{2}+a_{2} r^{4}+\ldots+a_{k} r^{2 k}+o\left(r^{2 k}\right)
$$

for any positive integer $k$ (and $a_{k}$ to be determined). From (3) and (4), we obtain

$$
\left(r^{\lambda}, \phi\right)=\Omega_{m} \int_{0}^{\infty} r^{\lambda+m-1} S_{\phi}(r) d r
$$

which indicates the application of $\Omega_{m} x_{+}^{\mu}$ with $\mu=\lambda+m-1$ to the testing function $S_{\phi}(r)$. Using the following Laurent series for $x_{+}^{\lambda}$ about $\lambda=-k$

$$
x_{+}^{\lambda}=\frac{(-1)^{k-1} \delta^{(k-1)}(x)}{(k-1)!(\lambda+k)}+x_{+}^{-k}+(\lambda+k) x_{+}^{-k} \ln x+\ldots
$$

we could show that the residue of $\left(r^{\lambda}, \phi(x)\right)$ at $\lambda=-m-2 k$ for nonnegative integer $k$ is given by

$$
\Omega_{m} \frac{\left(\delta^{(2 k)}, \phi(x)\right)}{(2 k)!}=\Omega_{m} \frac{S_{\phi}^{(2 k)}(0)}{(2 k)!}
$$

On the other hand, the residue of the function $r^{\lambda}$ of (2) for the same value of $\lambda$ is

$$
\frac{\Omega_{m} \Delta^{k} \delta(x)}{2^{k} k!m(m+2) \cdots(m+2 k-2)}
$$

Therefore we get

$$
S_{\phi}^{(2 k)}(0)=\frac{(2 k)!\Delta^{k} \phi(0)}{2^{k} k!m(m+2) \cdots(m+2 k-2)}
$$

This result can be used to write out the Taylor's series for $S_{\phi}(r)$, namely

$$
\begin{aligned}
S_{\phi}(r) & =\phi(0)+\frac{1}{2!} S_{\phi}^{\prime \prime}(0) r^{2}+\ldots+\frac{1}{(2 k)!} S_{\phi}^{(2 k)}(0) r^{2 k}+\ldots \\
& =\sum_{k=0}^{\infty} \frac{\Delta^{k} \phi(0) r^{2 k}}{2^{k} k!m(m+2) \cdots(m+2 k-2)}
\end{aligned}
$$

which is the well-known Pizetti's formula.

## 2. THE PRODUCT $\boldsymbol{r}^{\boldsymbol{- k}} \cdot \boldsymbol{\nabla} \boldsymbol{\delta}$

Let $\rho(x)$ be a fixed infinitely differentiable function defined on $R$ with the following properties:
(i) $\rho(x) \geq 0$,
(ii) $\rho(x)=0$ for $|x| \geq 1$,
(iii) $\rho(x)=\rho(-x)$,
(iv) $\int_{-1}^{1} \rho(x) d x=1$.

The function $\delta_{n}(x)$ is defined by $\delta_{n}(x)=n \rho(n x)$ for $n=1,2, \ldots$. It follows that $\left\{\delta_{n}(x)\right\}$ is a regular sequence of infinitely differentiable functions converging to the Dirac delta-function $\delta(x)$.

Now let $\mathcal{D}$ be the space of infinitely differentiable functions of a single variable with compact support and let $\mathcal{D}^{\prime}$ be the space of distributions defined on $\mathcal{D}$. Then if $f$ is an arbitrary distribution in $\mathcal{D}^{\prime}$, we define

$$
f_{n}(x)=\left(f * \delta_{n}\right)(x)=\left(f(t), \delta_{n}(x-t)\right) \quad \text { for } \quad n=1,2, \ldots
$$

It follows that $\left\{f_{n}(x)\right\}$ is a regular sequence of infinitely differentiable functions converging to the distribution $f(x)$ in $\mathcal{D}^{\prime}$.

The following definition for the non-commutative neutrix product $f \cdot g$ of two distributions $f$ and $g$ in $\mathcal{D}^{\prime}$ was given by Fisher in [9].

Definition 1. Let $f$ and $g$ be distributions in $\mathcal{D}^{\prime}$ and Let $g_{n}=g * \delta_{n}$. We say that the neutrix product $f \cdot g$ of $f$ and $g$ exists and is equal to $h$ if

$$
N-\lim _{n \rightarrow \infty}\left(f g_{n}, \phi\right)=(h, \phi)
$$

for all functions $\phi$ in $\mathcal{D}$, where $N$ is the neutrix (see [10]) having domain $N^{\prime}=\{1,2, \ldots\}$ and range $N^{\prime \prime}$ the real numbers, with negligible functions that are finite linear sums of the functions $n^{\lambda} \ln ^{r-1} n, \ln ^{r} n(\lambda>0, r=1,2, \ldots)$ and all functions of $n$ which converge to zero in the normal sense as $n$ tends to infinity.

The product of Definition 1 is not symmetric and hence $f \cdot g \neq g \cdot f$ in general.
In order to give a definition for a neutrix product $f \cdot g$ of two distributions in $\mathcal{D}_{m}^{\prime}$, the space of distributions defined $\mathcal{D}_{m}$. We may attempt to define a $\delta$-sequence in $\mathcal{D}_{m}$ by simply putting (see [12])

$$
\delta_{n}\left(x_{1}, \ldots, x_{m}\right)=\delta_{n}\left(x_{1}\right) \cdots \delta_{n}\left(x_{m}\right),
$$

where $\delta_{n}$ is defined as above. However, this definition is very difficult to use for distributions in $\mathcal{D}_{m}^{\prime}$ which are functions of $r$. We therefore consider the following approach (see [7]).

Let $\rho(s)$ be a fixed infinitely differentiable function defined on $R^{+}=[0, \infty)$ having the properties:
(i) $\rho(s) \geq 0$;
(ii) $\rho(s)=0$ for $s \geq 1$;
(iii) $\int_{R^{m}} \delta_{n}(x) d x=1$,
where $\delta_{n}(x)=c_{m} n^{m} \rho\left(n^{2} r^{2}\right)$ and $c_{m}$ is the constant satisfying (iii). It follows that $\left\{\delta_{n}(x)\right\}$ is a regular $\delta$-sequence of infinitely differentiable functions converging to $\delta(x)$ in $\mathcal{D}_{m}^{\prime}$.
Definition 2. Let $f$ and $g$ be distributions in $\mathcal{D}_{m}^{\prime}$ and let

$$
g_{n}(x)=\left(g * \delta_{n}\right)(x)=\left(g(x-t), \delta_{n}(t)\right), \quad \text { where } \quad t=\left(t_{1}, t_{2}, \ldots, t_{m}\right) .
$$

We say that the neutrix product $f \cdot g$ of $f$ and $g$ exists and is equal to $h$ if

$$
N-\lim _{n \rightarrow \infty}\left(f g_{n}, \phi\right)=(h, \phi),
$$

where $\phi \in \mathcal{D}_{m}$ and the $N$-limit is defined as above.
With Definition 2 and the normalization procedure of $\mu(x) x_{+}^{\lambda}, \mathrm{Li}$ (also in [7]) shows that the non-commutative neutrix product $r^{-k} \cdot \nabla \delta$, ( $\nabla$ is the gradient operator, that is, $\left.\nabla=\sum_{i=1}^{m} \frac{\partial}{\partial x_{i}}\right)$ exists and

$$
\begin{equation*}
r^{-2 k} \cdot \nabla \delta=-\frac{1}{2^{k+1}(k+1)!(m+2) \cdots(m+2 k)} \sum_{i=1}^{m}\left(x_{i} \Delta^{k+1} \delta\right) \tag{5}
\end{equation*}
$$

and

$$
\begin{equation*}
r^{1-2 k} \cdot \nabla \delta=0 \tag{6}
\end{equation*}
$$

where $k$ is a positive integer.

## 3. THE DISTRIBUTION $\Delta r^{2-m}$

We note that $r^{2-m}$ for $m \geq 3$ is a harmonic function in any region which does not contain the origin, i.e., that $\Delta r^{2-m}$ vanishes in the ordinary sense for all $r \neq 0$. For the case of generalized functions we have

$$
\left(\Delta r^{2-m}, \phi\right)=\left(r^{2-m}, \Delta \phi\right)=\lim _{e \rightarrow 0^{+}} \int_{r \geq \epsilon} \frac{\Delta \phi}{r^{m-2}} d v
$$

Now applying Green's theorem to this integral, we have

$$
\int_{r \geq \epsilon} \frac{\Delta \phi}{r^{m-2}} d v=\int_{r \geq \epsilon} \phi \Delta r^{2-m} d v-\int_{r=\epsilon} \frac{\partial \phi}{\partial r} r^{2-m} d s+\int_{r=\epsilon} \phi \frac{\partial}{\partial r} r^{2-m} d s
$$

where $d s$ is the element of area on the sphere $r=\epsilon$. Now

$$
\int_{r \geq e} \phi \Delta r^{2-m} d v=0
$$

since outside the ball $r<\epsilon$ the function $r^{2-m}$ is harmonic. As for the second term,

$$
\int_{r=\epsilon} \frac{\partial \phi}{\partial r} r^{2-m} d s=\epsilon^{2-m} \int_{r=\epsilon} \frac{\partial \phi}{\partial r} d s=O(\epsilon)
$$

by noting that $\partial \phi / \partial r$ is bounded near the origin. On the other hand,

$$
\int_{r=\epsilon} \phi \frac{\partial}{\partial r} r^{2 \sim m} d s=(2-m) \epsilon^{1-m} \int_{r=\epsilon} \phi d s=(2-m) \Omega_{m} S_{\varepsilon}(\phi)
$$

where $S_{e}(\phi)$ is the mean value of $\phi(x)$ on the sphere of radius $\epsilon$. In the limit as $\epsilon \rightarrow 0^{+}$, of course, $S_{\varepsilon}(\phi) \rightarrow \phi(0)$, so that

$$
\left(\Delta r^{2-m}, \phi\right)=\lim _{e \rightarrow 0^{+}} \int_{r \geq \epsilon} \frac{\Delta \phi}{r^{m-2}} d v=(2-m) \Omega_{m} \phi(0)
$$

Hence we may write

$$
\begin{equation*}
\frac{1}{(2-m) \Omega_{m}} \Delta r^{2-m}=\delta(x) \tag{7}
\end{equation*}
$$

and it immediately follows that

$$
E=\frac{\Gamma(m / 2)}{(2-m) 2 \pi^{m / 2}} r^{2-m}
$$

is an elementary solution of Laplacian equation $\Delta u(x)=\delta(x)$.
A similar calculation for dimension $m=2$ leads to the result

$$
\begin{equation*}
\frac{1}{2 \pi} \Delta \ln r=\delta(x) \tag{8}
\end{equation*}
$$

## 4. SOME PRODUCTS

In this section, we will provide a nicer version of the main theorem in [7] as well as the results obtained in [6] to provide the products outlined in the abstract.

It obviously follows from equations (5) and (7) that
Theorem 1. The non-commutative neutrix product $r^{-k} \cdot \nabla\left(\Delta r^{2-m}\right)$ exists for $m \geq 3$. Furthermore

$$
\begin{equation*}
r^{-2 k} \cdot \nabla\left(\Delta r^{2-m}\right)=\frac{(m-2) \pi^{m / 2}}{2^{k}(k+1)!(m+2) \cdots(m+2 k) \Gamma(m / 2)} \sum_{i=1}^{m}\left(x_{i} \Delta^{k+1} \delta\right) \tag{9}
\end{equation*}
$$

and

$$
\begin{equation*}
r^{1-2 k} \cdot \nabla\left(\Delta r^{2-m}\right)=0 \tag{10}
\end{equation*}
$$

where $k$ is a positive integer.
Remark. The product of $x_{i}$ and $\Delta^{k+1} \delta$ in our theorem is well defined since

$$
\begin{equation*}
\left(x_{i} \Delta^{k+1} \delta, \phi\right)=\left(\delta, \Delta^{k+1}\left(x_{i} \phi\right)\right) \tag{11}
\end{equation*}
$$

The following lemma will play an important role in simplifying the above theorem.

## Lemma 1.

$$
\begin{equation*}
\sum_{i=1}^{m} \Delta^{k+1}\left(x_{i} \phi\right)=2(k+1) \nabla\left(\Delta^{k} \phi\right)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{k+1} \phi \quad \text { for } \quad k \geq 0 \tag{12}
\end{equation*}
$$

Proof. We use an inductive method to show the lemma. It is obviously true for $k=0$. When $k=1$, we have $\Delta^{2}\left(x_{i} \phi\right)=4 \frac{\partial}{\partial x_{i}} \Delta \phi+x_{i} \Delta^{2} \phi$ simply by calculating the left-hand side. Then,

$$
\sum_{i=1}^{m} \Delta^{2}\left(x_{i} \phi\right)=4 \nabla(\Delta \phi)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{2} \phi
$$

Assume, equation (12) holds for the case of $k-1$, that is

$$
\sum_{i=1}^{m} \Delta^{k}\left(x_{i} \phi\right)=2 k \nabla\left(\Delta^{k-1} \phi\right)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{k} \phi
$$

Hence it follows that

$$
\begin{aligned}
& \sum_{i=1}^{m} \Delta^{k+1}\left(x_{i} \phi\right)=\Delta \sum_{i=1}^{m} \Delta^{k}\left(x_{i} \phi\right)=\Delta\left\{2 k \nabla\left(\Delta^{k-1} \phi\right)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{k} \phi\right\} \\
& \quad=2 k \nabla\left(\Delta^{k} \phi\right)+\sum_{i=1}^{m} \Delta\left(x_{i} \Delta^{k} \phi\right)=2 k \nabla\left(\Delta^{k} \phi\right)+\sum_{i=1}^{m}\left\{2 \frac{\partial}{\partial x_{i}} \Delta^{k} \phi+x_{i} \Delta^{k+1} \phi\right\} \\
& =2 k \nabla\left(\Delta^{k} \phi\right)+2 \nabla\left(\Delta^{k} \phi\right)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{k+1} \phi=2(k+1) \nabla\left(\Delta^{k} \phi\right)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{k+1} \phi
\end{aligned}
$$

This completes the proof of Lemma 1.

Lemma 2. $\quad \sum_{i=1}^{m}\left(x_{i} \Delta^{k+1} \delta\right)=-2(k+1) \nabla\left(\Delta^{k} \delta\right), \quad$ where $\quad k \geq 0$.
Proof. Applying equation (12), we have

$$
\begin{aligned}
\left(\sum_{i=1}^{m}\left(x_{i} \Delta^{k+1} \delta\right), \phi\right) & =\left(\delta, \sum_{i=1}^{m} \Delta^{k+1}\left(x_{i} \phi\right)\right) \\
& =\left(\delta,\left\{2(k+1) \nabla\left(\Delta^{k} \phi\right)+\left(\sum_{i=1}^{m} x_{i}\right) \Delta^{k+1} \phi\right\}\right) \\
& =2(k+1) \nabla\left(\Delta^{k} \phi(0)\right)=\left(-2(k+1) \nabla\left(\Delta^{k} \delta\right), \phi\right)
\end{aligned}
$$

Therefore, we have reached our conclusion in Lemma 2.
Now, using equations (9) and (13), we can prove the following theorem.
Theorem 2. The non-commutative neutrix product $r^{-k} \cdot \nabla\left(\Delta r^{2-m}\right)$ exists for $m \geq 3$. Further

$$
\begin{equation*}
r^{-2 k} \cdot \nabla\left(\Delta r^{2-m}\right)=\frac{(2-m) \pi^{m / 2}}{2^{k-1} k!(m+2) \cdots(m+2 k) \Gamma(m / 2)} \nabla\left(\Delta^{k} \delta\right) \tag{14}
\end{equation*}
$$

and

$$
\begin{equation*}
r^{1-2 k} \cdot \nabla\left(\Delta r^{2-m}\right)=0 \tag{15}
\end{equation*}
$$

where $k$ is a positive integer.
In particular, we obtain

$$
\begin{aligned}
\left(x_{1}^{2}+x_{2}^{2}+x_{3}^{2}\right)^{-1} \cdot \nabla\left(\Delta\left(x_{1}^{2}+x_{2}^{2}+x_{3}^{2}\right)^{-1 / 2}\right) & =-\frac{2 \pi}{5} \nabla(\Delta \delta) \\
\left(x_{1}^{2}+x_{2}^{2}+x_{3}^{2}\right)^{-1 / 2} \cdot \nabla\left(\Delta\left(x_{1}^{2}+x_{2}^{2}+x_{3}^{2}\right)^{-1 / 2}\right) & =0 \\
\left(x_{1}^{2}+x_{2}^{2}+x_{3}^{2}+x_{4}^{2}\right)^{-2} \cdot \nabla\left(\Delta\left(x_{1}^{2}+x_{2}^{2}+x_{3}^{2}+x_{4}^{2}\right)^{-1}\right) & =-\frac{\pi^{2}}{96} \nabla\left(\Delta^{2} \delta\right)
\end{aligned}
$$

by noting that $\Gamma(3 / 2)=\sqrt{\pi} / 2$.
The following result can be found in [6].
Theorem 3. The non-commutative neutrix product $r^{-k} \cdot \delta$ exists. Furthermore

$$
\begin{equation*}
r^{-2 k} \cdot \delta=\frac{\Delta^{k} \delta}{2^{k} k!m(m+2) \cdots(m+2 k-2)} \tag{16}
\end{equation*}
$$

and

$$
\begin{equation*}
r^{1-2 k} \cdot \delta=0 \tag{17}
\end{equation*}
$$

where $k$ is a positive integer.
Theorem 4. The non-commutative neutrix products $r^{-k} \cdot \Delta \ln r$ and $r^{-k} \cdot \nabla(\Delta \ln r)$ exist for $m=2$, and

$$
r^{-2 k} \cdot \Delta \ln r=\frac{\pi \Delta^{k} \delta}{2^{2 k-1}(k!)^{2}}, \quad r^{-2 k} \cdot \nabla(\Delta \ln r)=\frac{\pi \nabla\left(\Delta^{k} \delta\right)}{2^{2 k-1} k!(k+1)!}
$$

where $k$ is a positive integer.

Proof. It immediately follows from equations (5), (8), (13) and (16).
In particular for $k=1$, we get

$$
\begin{gathered}
\left(x_{1}^{2}+x_{2}^{2}\right)^{-1} \cdot \Delta \ln \left(x_{1}^{2}+x_{2}^{2}\right)^{1 / 2}=\frac{\pi \Delta \delta}{2} \\
\left(x_{1}^{2}+x_{2}^{2}\right)^{-1} \cdot \nabla\left(\Delta \ln \left(x_{1}^{2}+x_{2}^{2}\right)^{12}\right)=\frac{\pi \nabla(\Delta \delta)}{4}
\end{gathered}
$$

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